

# Cytochrome $c_{552}$ Mutants: Structure and Dynamics at the Active Site Probed by Multidimensional NMR and Vibration Echo Spectroscopy<sup>†</sup>

Aaron M. Massari,<sup>‡</sup> Brian L. McClain,<sup>‡,||</sup> Ilya J. Finkelstein,<sup>‡</sup> Andrew P. Lee,<sup>§</sup> Heather L. Reynolds,<sup>§</sup> Kara L. Bren,<sup>§</sup> and Michael D. Fayer<sup>\*,‡</sup>

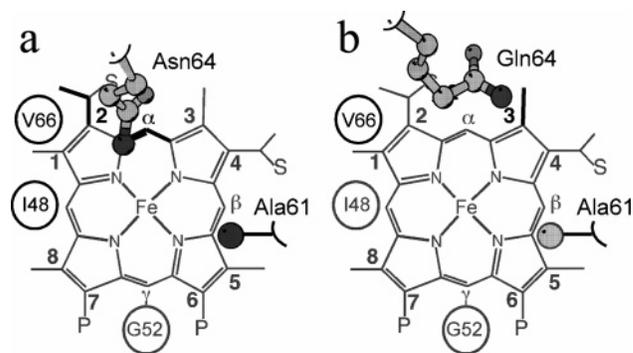
Department of Chemistry, Stanford University, Stanford, California 94305, and  
Department of Chemistry, University of Rochester, Rochester, New York 14627-0216

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Spectrally resolved infrared stimulated vibrational echo experiments are used to measure the vibrational dephasing of a CO ligand bound to the heme cofactor in two mutated forms of the cytochrome  $c_{552}$  from *Hydrogenobacter thermophilus*. The first mutant (*Ht*-M61A) is characterized by a single mutation of Met61 to an Ala (*Ht*-M61A), while the second variant is doubly modified to have Gln64 replaced by an Asn in addition to the M61A mutation (*Ht*-M61A/Q64N). Multidimensional NMR experiments determined that the geometry of residue 64 in the two mutants is consistent with a non-hydrogen-bonding and hydrogen-bonding interaction with the CO ligand for *Ht*-M61A and *Ht*-M61A/Q64N, respectively. The vibrational echo experiments reveal that the shortest time scale vibrational dephasing of the CO is faster in the *Ht*-M61A/Q64N mutant than that in *Ht*-M61A. Longer time scale dynamics, measured as spectral diffusion, are unchanged by the Q64N modification. Frequency–frequency correlation functions (FFCFs) of the CO are extracted from the vibrational echo data to confirm that the dynamical difference induced by the Q64N mutation is primarily an increase in the fast (hundreds of femtoseconds) frequency fluctuations, while the slower (tens of picoseconds) dynamics are nearly unaffected. We conclude that the faster dynamics in *Ht*-M61A/Q64N are due to the location of Asn64, which is a hydrogen bond donor, above the heme-bound CO. A similar difference in CO ligand dynamics has been observed in the comparison of the CO derivative of myoglobin (MbCO) and its H64V variant, which is caused by the difference in axial residue interactions with the CO ligand. The results suggest a general trend for rapid ligand vibrational dynamics in the presence of a hydrogen bond donor.

## I. Introduction

Cytochromes *c* (cyt *c*'s) are small electron-transfer heme proteins that play key roles in respiration and photosynthesis.<sup>1–3</sup> Due to their wide availability, robust structure, and ease of handling, these proteins have served as model systems to study the interplay of protein structure, dynamics, and function.<sup>4–13</sup> Cytochrome  $c_{552}$  from *Hydrogenobacter thermophilus* (*Ht*-cyt  $c_{552}$ )<sup>14–17</sup> was recently shown to be a structurally unique example within the cyt  $c_8$  family of class I<sup>18</sup> cyt *c*'s. Proteins in the cyt  $c_8$  structural family, for example, *Pseudomonas* cyt  $c_{551}$ 's,<sup>19–21</sup> typically have an asparagine residue at position 64 (Asn64) that is situated to donate a hydrogen bond to a methionine at position 61 (Met61), which occupies the sixth coordination site to the heme (Figure 1a; sequence numbering is based on the *Pseudomonas aeruginosa* cyt  $c_{551}$  sequence).<sup>19–21</sup> This interaction has been shown to play an important role in determining the axial Met orientation<sup>22,23</sup> and heme redox potential.<sup>24</sup> Unlike its homologues, *Ht*-cyt  $c_{552}$  has a glutamine at position 64 (Gln64), which is not oriented toward the axial Met61, but is instead localized over a methyl group (heme 3-CH<sub>3</sub> using the Fischer numbering system) at the heme edge, near the protein surface (Figure 1b).<sup>15–17</sup> Zhong and co-workers have shown that the absence



**Figure 1.** Schematic representation of the heme and its environment in (a) *Ht*-M61A/Q64N and (b) *Ht*-M61A. The heme axial ligands (His/CO) are omitted for clarity. The approximate positions of Ile48, Gly52, and Val66 are shown with circles, and Ala61 and residue 64 are shown in ball-and-stick representations. All amino acids shown are oriented above the heme plane (toward the viewer) as is the axial CO ligand. Amino acids and heme substituents that show NOEs with the side-chain NH<sub>2</sub> group of residue 64 are highlighted in black.

of a hydrogen bond to the axial Met affects the dynamics, and presumably the functionality, of the active site by imparting a fluxional character to Met61 and influencing heme redox potential.<sup>17,22–24</sup>

Multidimensional spectroscopic techniques have provided a wealth of information in the field of biology. Two-dimensional NMR studies have elucidated the structural and dynamical details of solvated proteins on time scales longer than tens of

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\* Author to whom correspondence should be addressed. E-mail: fayer@stanford.edu.

<sup>‡</sup> Stanford University.

<sup>§</sup> University of Rochester.

<sup>||</sup> Permanent address: Department of Chemistry and Biochemistry, California State University at Long Beach, Long Beach, CA 90840.

picoseconds.<sup>13,25–33</sup> Multidimensional IR spectroscopy has improved upon the dynamic range of NMR techniques, revealing biochemical dynamics that occur on the picosecond to femtosecond time scales.<sup>34–41</sup> Vibrational echo spectroscopy is a multidimensional IR technique that is sensitive to the relationship between structure and dynamics in heme proteins.<sup>34–43</sup> These measurements probe the dephasing of a heme-bound CO vibration caused by structural fluctuation of the protein. There is also a contribution to the vibrational echo observable from vibrational energy relaxation from the CO vibration to other modes. In carbonmonoxy heme proteins, such as myoglobin (MbCO), hemoglobin (HbCO), and cyt c, the vibrational lifetime of the CO stretch is sufficiently long that the vibrational echo decay is caused primarily by protein structural fluctuations. Within the electrostatic force model described below,<sup>39,44</sup> the protein environment is envisioned as a network of partial charges, whose movement generates a time-dependent electric field that influences the CO vibrational frequency to produce dynamic dephasing.<sup>34,35,45–48</sup> Nonlinear response theory<sup>49</sup> allows the extraction of the equilibrium autocorrelation function of the fluctuations in the CO vibrational frequency, or frequency–frequency correlation function (FFCF). The FFCF provides a quantitative description of the dynamics measured in vibrational echo spectroscopy that is useful for comparing the dynamics of different proteins.

In the current work, spectrally resolved infrared stimulated (three-pulse) vibrational echo spectroscopy<sup>36</sup> was used to directly measure the active site structural dynamics of two mutated forms of *Ht*-cyt c<sub>552</sub>.<sup>17,22,50</sup> In both mutants, the heme axial ligand, Met61, was replaced by an alanine (M61A), which allowed a strong IR probe (CO) to be strategically bound to the ferrous heme in place of Met61 to report on the active site dynamics of both proteins. The first mutant studied (*Ht*-M61A) was characterized by only the single M61A mutation and retained an axial Gln64, which was expected to be oriented out of the heme pocket (Figure 1b) as it is in the native protein.<sup>16,17,22</sup> The second mutant (*Ht*-M61A/Q64N) was further modified to have Gln64 replaced by an Asn in an effort to generate an active site structure similar to that seen in most cyt c<sub>8</sub>'s,<sup>19–23</sup> with Asn64 positioned to interact with the heme-bound CO (Figure 1a).

The vibrational echo data revealed that the ultrafast active site dynamics ( $\leq 1$  ps) sensed by the heme-bound CO in *Ht*-M61A/Q64N are noticeably faster than those for *Ht*-M61A, while the rates of slower processes (tens of picoseconds time scale) are very similar for the two variants. Multidimensional NMR determined that the geometry of residue 64 in the two mutants produced conformations consistent with a non-hydrogen-bonding and hydrogen-bonding interaction between residue 64 and the CO ligand for *Ht*-M61A and *Ht*-M61A/Q64N, respectively. Gln64 was found to localize away from the heme pocket in *Ht*-M61A, as it does in the native *Ht*-cyt c, while Asn64 is oriented into the active site in *Ht*-M61A/Q64N to donate a hydrogen bond as it would in typical cyt c<sub>8</sub>'s. The faster dynamics measured in *Ht*-M61A/Q64N are attributed to the interaction of the hydrogen-bond-donating Asn64 with the heme-bound CO. The dynamical changes observed in this engineered ligand-binding heme protein are analogous to those reported previously for the natural ligand-binding heme protein MbCO. A variant of MbCO in which the hydrogen-bond-donating distal histidine (coincidentally located at position 64) is replaced with a valine (H64V) displays slower vibrational dephasing than the native protein.<sup>51</sup> These examples suggest a general trend toward rapid active site dynamics, as sensed by the CO ligand vibrational dephasing, in the presence of a hydrogen bond donor,

and may represent a mechanism by which an organism imparts a unique selectivity to substrate binding or reactivity at the active site of a protein or enzyme.

## II. Materials and Methods

**A. Protein Expression and Purification.** Preparation of *Ht*-M61A and *Ht*-M61A/Q64N utilized an *E. coli*-based expression system.<sup>50,52</sup> Molecular biology procedures and materials and the preparation of *Ht*-M61A are described in detail elsewhere.<sup>22,53</sup> To prepare *Ht*-M61A/Q64N, the polymerase chain reaction overlap extension method<sup>54</sup> was employed using the pSHC552A61 expression plasmid<sup>53</sup> as the template. The mutagenic primers (mutation site italicized) were 5'-CCC GCGCTCCTAATAATGTAACC-3' and 5'-CGGTTCATTATTAGGAGGCGCGG-3'. Cloning, expression, and purification of *Ht*-M61A/Q64N were as described for *Ht*-M61A.<sup>53</sup>

**B. Sample Preparation for Vibrational Echo Spectroscopy.** To prepare aqueous samples of carbonmonoxy *Ht*-M61A and *Ht*-M61A/Q64N, 10 mg of lyophilized protein was dissolved in 1.0 mL of pH 7.0 H<sub>2</sub>O phosphate buffer (50 mM). The buffer pH was measured before addition of protein. The solutions were reduced with a 5-fold excess of dithionite (Sigma Aldrich) and stirred under a CO atmosphere for 1 h. The solutions were centrifuged at 3000 relative centrifugal force for 15 min through a 0.45  $\mu$ m acetate filter (Pall Nanosep MF) to remove particulates. The samples were further concentrated by repeated centrifugation (Eppendorf 5415D) over modified polyethersulfone membranes (Pall Nanosep 3K Omega) to a final protein concentration of 10–15 mM. The sample was then placed in a sample cell with CaF<sub>2</sub> windows and a 50  $\mu$ m Teflon spacer. UV–visible (Varian Cary 3E) and Fourier transform infrared (FTIR; ATI Mattson Infinity 9495) absorption spectroscopies were performed to determine all protein concentrations. The samples had mid-IR absorbances of 0.1 at the CO stretching frequency on a background absorbance of 0.5.

**C. NMR Spectroscopy.** Samples of carbonmonoxy *Ht*-M61A and *Ht*-M61A/Q64N for NMR (3–4 mM protein in 50 mM sodium phosphate, 10% D<sub>2</sub>O, pH 4) were prepared by flushing samples with CO, followed by introducing a  $\sim 20$ -fold molar excess of Na<sub>2</sub>S<sub>2</sub>O<sub>4</sub>. <sup>1</sup>H NMR data were collected on a Varian INOVA 500-MHz spectrometer at 298 K. Two-dimensional nuclear Overhauser effect (NOE) spectra were collected with 4096 points in the F2 dimension, 512 increments in the F1 dimension, and a 12 000-Hz spectral width. The mixing time was 100 ms, and the recycle time was 1.3 s. Solvent suppression was achieved by presaturation.

NMR data processing and analysis were performed using FELIX 97 (Accelrys). Assignments of selected <sup>1</sup>H resonances of *Ht*-M61A and *Ht*-M61A/Q64N were made by standard methods<sup>55</sup> aided by comparison to published <sup>1</sup>H resonance assignments for reduced *Ht* cyt c<sub>552</sub><sup>15,50</sup> and *Ht*-Q64N.<sup>22</sup>

**D. Stimulated Vibrational Echo Spectroscopy.** The experimental setup has been previously described in detail.<sup>56</sup> Briefly, tunable mid-IR pulses with a center frequency adjusted to match the center frequency of the protein sample of interest (1965–1976 cm<sup>-1</sup>) were generated by an optical parametric amplifier pumped with a regeneratively amplified Ti:sapphire laser. The bandwidth and pulse duration used in these experiments were 150 cm<sup>-1</sup> and 100 fs, respectively. The mid-IR pulse was split into three temporally controlled pulses ( $\sim 700$  nJ/pulse). The delay between the first two pulses,  $\tau$ , was scanned at each time  $T_w$ , the delay between pulses two and three. The three beams were crossed and focused at the sample. The spot size at the sample was  $\sim 150$   $\mu$ m. The vibrational echo pulse

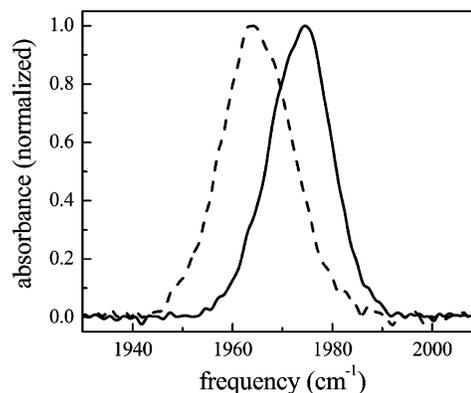
generated in the phase-matched direction was dispersed through a 0.5 m monochromator (1.2 cm<sup>-1</sup> spectral resolution) and detected with either a liquid-nitrogen-cooled HgCdTe array detector (Infrared Associates/Infrared Systems Development) or a liquid-nitrogen-cooled InSb single element detector (EG&G Judson). A power dependence study was performed on all samples, and the data showed no power-dependent effects.<sup>57</sup> Data collection for all samples was performed at room temperature in an enclosed, dry-air-purged environment.

**E. FFCF Extraction from Vibrational Echo Data.** To extract quantitative information from the vibrational echo data, nonlinear response theory calculations were compared to the experimental data.<sup>39,49</sup> Within conventional approximations,<sup>49</sup> both the vibrational echo and the linear-infrared absorption spectrum are completely determined by the FFCF. The FFCF is the starting point from which the experimental observables are calculated. A multiexponential form of the FFCF,  $C(t)$ , was used in accord with previous vibrational echo analysis and molecular dynamics simulations of heme proteins.<sup>39,51,53</sup> It is important to note that exponentials in the FFCF do not give rise to single- or multiple-exponential vibrational echo decay curves but rather produce complicated nonexponential decays. The FFCF has the form

$$C(t) = \Delta_0^2 + \sum_{i=1}^n \Delta_i^2 \exp(-t/\tau_i) \quad (1)$$

Here,  $\Delta_0$  is the contribution from inhomogeneous broadening. Inhomogeneous broadening is caused by variations in protein structure that influence the CO frequency but evolve on time scales that are much slower than the experimental time window. In this study, the structural dynamics that occur on the time scale of  $\sim 100$  ps will contribute to inhomogeneous broadening.  $\Delta_i$  is the magnitude of the contribution from a frequency-perturbing process with correlation time  $\tau_i$ . If  $\tau_i$  is fast compared to  $\Delta_i^{-1}$  ( $\Delta_i \tau_i \ll 1$ ,  $\Delta$  in radians/ps) for a given exponential term, then that component of the FFCF is motionally narrowed.<sup>58–61</sup> For a motionally narrowed term in  $C(t)$ ,  $\Delta$  and  $\tau_i$  cannot be determined independently,<sup>53</sup> but a pure dephasing time,  $T_2^*$ , can be defined ( $T_2^* = (\Delta^2 \tau)^{-1}$ ), which describes the “homogeneous line width” for that component of the FFCF. Although protein dynamics generally occur over a continuum of time scales, a multiexponential  $C(t)$  organizes these fluctuations into experimentally relevant time scales that can be compared from one system to another.

The FFCF is used to calculate the linear absorption spectrum and a series of vibrational echo decay curves ( $\tau$  scanned,  $T_w$  fixed) for a range of  $T_w$  values. An FFCF with a single exponential plus a constant term was not able to simultaneously reproduce the linear-IR spectrum and the vibrational echo decay curves at all  $T_w$  values for either protein in this study. Two distinct FFCFs for *Ht-M61A* and *Ht-M61A/Q64N*, composed of a biexponential plus a constant term, were adequate to simultaneously fit the absorption spectrum and the vibrational echo decay curves of each protein. To maximize the efficiency of the empirical fits,  $\delta$ -function laser pulses were used when fitting the data. A comparison of these fits to those obtained by performing the full three time-ordered integrals with the finite experimental pulse durations<sup>49</sup> verified that the effects of pulse duration were negligibly small given the very short pulses used in the experiments. The FFCF obtained from analysis of the data using response theory calculations was deemed correct when it could be used to calculate vibrational echo decays that



**Figure 2.** Normalized FTIR spectra of the CO stretching mode bound to *Ht-M61A* (solid line) and *Ht-M61A/Q64N* (dashed line).

fit the experimental vibrational echo data at all  $T_w$  values and simultaneously reproduce the linear absorption spectrum.

### III. Results and Discussion

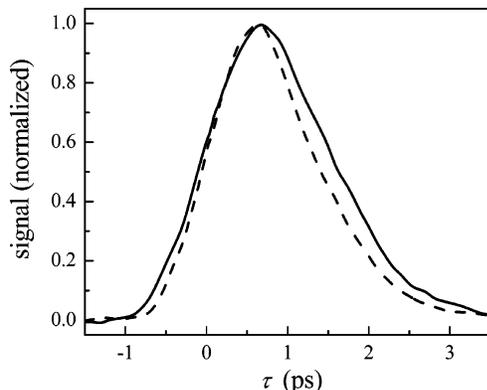
**A. Linear-IR Spectroscopy.** The background-subtracted linear FTIR spectra of CO bound to *Ht-M61A* and *Ht-M61A/Q64N* are shown in Figure 2. All peaks have been fit as Gaussian distributions to determine their full width at half-maxima (fwhm's) and center frequencies (Table 1). The spectrum of *Ht-M61A* (solid line) shows a single transition at 1974 cm<sup>-1</sup> with a fwhm of 14.7 cm<sup>-1</sup>, while *Ht-M61A/Q64N* (dashed line) exhibits a single transition at 1965 cm<sup>-1</sup> with a fwhm of 16.8 cm<sup>-1</sup>. The red shifting of the CO stretching frequency of *Ht-M61A/Q64N* relative to *Ht-M61A* is reminiscent of the IR spectrum of aqueous MbCO, which exhibits three primary CO stretching peaks corresponding to three structurally distinct conformational substates.<sup>39,51,56,62–66</sup> The A<sub>0</sub> band (1965 cm<sup>-1</sup>) is attributed to a conformation in which the distal histidine (His64) is positioned out of the heme pocket,<sup>51,67–71</sup> while the A<sub>1</sub> (1944 cm<sup>-1</sup>) and A<sub>3</sub> (1938 cm<sup>-1</sup>) bands arise from the distal histidine being localized in the heme pocket in two distinct orientational geometries.<sup>39,72</sup> In the case of MbCO, the inclusion of a polar, hydrogen-bond-donating residue inside the heme pocket results in a spectral shift of the CO stretch to lower frequencies. When this example is compared to the spectra of *Ht-M61A* and *Ht-M61A/Q64N* in Figure 2, the observed red spectral shift caused by replacing Gln64 with an Asn implies that the Gln may be located out of the active site (as does His64 in the MbCO A<sub>0</sub> substate) whereas the Asn is directed into the heme pocket (MbCO A<sub>1</sub> or A<sub>3</sub> substate). According to the electrostatic force model,<sup>39,44</sup> the motions of a charged moiety in the vicinity of the heme-bound CO should dramatically affect the vibrational dynamics of this chromophore. However, the Gaussian shape of the spectral bands in *Ht-M61A* and *Ht-M61A/Q64N* suggests that these transitions are inhomogeneously broadened, which would obscure any dynamical information contained in the linear spectra. The vibrational stimulated echo experiments described below confirm that the spectral bands are indeed inhomogeneously broadened, and therefore, such experiments are necessary to uncover the underlying dynamics.

**B. Multidimensional IR Studies: Stimulated Vibrational Echo Spectroscopy.** Figure 3 shows the vibrational echo decays for *Ht-M61A* (solid curve) and *Ht-M61A/Q64N* (dashed curve) at  $T_w = 0.5$  ps. The decay of *Ht-M61A/Q64N* is noticeably faster than that of *Ht-M61A*, indicating that the CO dephases faster when Gln64 is replaced by Asn64. For clarity, we focus our attention here on a single  $T_w$ ; however, this trend is consistent at all values of  $T_w$ . In vibrational echo experiments,

**TABLE 1: FTIR Peak Centers and Line Widths and the Best-Fit  $C(t)$  Parameters for *Ht-M61A* and *Ht-M61A/Q64N*<sup>a</sup>**

protein	FTIR peak (cm <sup>-1</sup> )	fwhm (cm <sup>-1</sup> )	$\Delta_0$ (cm <sup>-1</sup> )	$\Delta_1$ (cm <sup>-1</sup> )	$\tau_1$ (ps)	$T_2^*$ (ps) (= $1/(\Delta_1^2\tau_1)$ )	$\Delta_2$ (cm <sup>-1</sup> )	$\tau_2$ (ps)
<i>Ht-M61A</i>	1974	14.7	4.49	4.73	0.31	4.1	3.0	5.3
<i>Ht-M61A/Q64N</i>	1965	16.8	5.19	4.85	0.41	2.9	3.8	8.7

<sup>a</sup> The pure dephasing times ( $T_2^*$ ) of the motionally narrowed component of the FFCFs are given since  $\Delta_1$  and  $\tau_1$  cannot be independently determined.

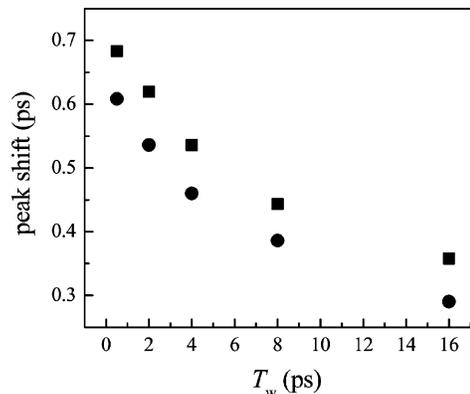


**Figure 3.** Spectrally resolved vibrational echo decays at  $T_w = 0.5$  ps for CO bound to *Ht-M61A* (1975 cm<sup>-1</sup>, solid line) and *Ht-M61A/Q64N* (1965 cm<sup>-1</sup>, dashed line).

a faster rate of dephasing indicates that the frequency of the heme-bound CO is fluctuating more rapidly. Within the electrostatic force model described above, structural fluctuations produce motions of charged and polar residues that induce the largest changes in CO transition frequency. The dynamic response of these two proteins as measured by vibrational echo spectroscopy confirms the implications of the linear-IR spectra discussed above. The increased fluctuations of the CO frequency are consistent with the structural picture in which Asn64 in *Ht-M61A/Q64N* is localized near the CO, whereas the Gln64 in *Ht-M61A* is directed out of the active site.

In stimulated (three-pulse) vibrational echo experiments, the dynamics that occur on time scales longer than those shown for a single vibrational echo decay (Figure 3) can be measured by varying the time delay between the second and third pulses,  $T_w$ . Although at each  $T_w$  the entire decay curve is measured, these dynamics, termed spectral diffusion, are conveniently depicted by plotting the vibrational echo peak shift<sup>73–76</sup> as a function of  $T_w$ . The vibrational echo peak shift is the difference between the time of peak amplitude of the decay curve and zero time. When longer time scale protein dynamics are present, measured in a vibrational echo experiment as frequency fluctuations that result in CO dephasing, the echo decays become faster as  $T_w$  becomes longer and their peaks shift toward the origin. In the frequency domain (Fourier transform of the vibrational echo decay), the changes observed in the vibrational echo decays with  $T_w$  show that the dynamical line width broadens with increasing  $T_w$  due to protein dynamics that influence the CO frequency on the  $T_w$  time scale. A broader dynamical line width (faster vibrational echo decay and smaller peak shift) means that a larger portion of the total possible structural configurations of the protein has been sampled. For long enough  $T_w$ , spectral diffusion is complete, and all chromophores have sampled the entire spectral line. In this case, the dynamic line shape is equal to the absorption line, and the vibrational echo peak shift is zero.

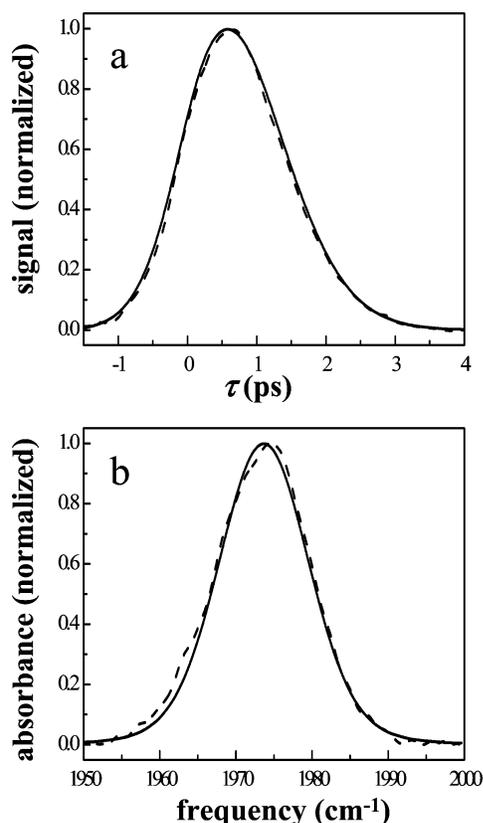
The vibrational echo peak shifts for *Ht-M61A* and *Ht-M61A/Q64N* are shown in Figure 4 as a function of  $T_w$ . The peak



**Figure 4.** Vibrational echo peak shifts as a function of  $T_w$  for *Ht-M61A* (filled squares) and *Ht-M61A/Q64N* (filled circles).

shift values for *Ht-M61A/Q64N* (filled circles) are consistently smaller than those for *Ht-M61A* (filled squares). This shows that for each  $T_w$  time delay, the CO bound to *Ht-M61A/Q64N* has sampled a greater fraction of the spectral line, which is consistent with the faster dephasing shown at a single  $T_w$  in Figure 3. The fact that the vibrational echo decay peaks have not shifted to zero by  $T_w = 16$  ps demonstrates that the full range of protein dynamics affecting the CO frequency have not occurred within this time frame; not all protein configurations that influence the frequency of the CO vibrational transition have been accessed. In these experiments, fluctuations on times greater than  $\sim 50$  ps appear as inhomogeneous broadening, which is accounted for by the  $\Delta_0$  term in eq 1. Aside from the nearly constant offset between the data for these two mutants (Figure 4), the peak shifts for both proteins as a function of  $T_w$  have the same qualitative shape. The protein dynamics, as sensed by the heme-bound CO, that occur on time scales longer than a few picoseconds qualitatively appear to be unaffected by replacing the Gln64 in *Ht-M61A* with the Asn64 in *Ht-M61A/Q64N*. This indicates that the primary influence of Asn64 on the CO dynamics occurs on the time scale of a few picoseconds or faster.

To compare the structural dynamics of these two proteins quantitatively, the FFCF for each protein was obtained by simultaneous fits that reproduce the linear-IR absorption spectrum and the vibrational echo decay curves at all  $T_w$  values using the procedure described in section II.E. The parameters for the best-fit FFCFs ( $C(t)$ ) for *Ht-M61A* and *Ht-M61A/Q64N* are summarized in Table 1. As explained in section II.E, the constant term ( $\Delta_0$ ) in  $C(t)$  accounts for static or quasi-static frequency distributions whose frequency fluctuations are very slow on the time scale of the experiments. An example of the quality of the fits obtained using only the five adjustable parameters in the biexponential FFCFs is shown for the vibrational echo decay at  $T_w = 2$  ps and the linear-IR spectrum for *Ht-M61A* in Figure 5. The overlaid fits to the experimental vibrational echo decay curves at all five  $T_w$  values and the linear-IR spectra calculated for both proteins using the parameters listed in Table 1 are available in the Supporting Information.



**Figure 5.** (a) Experimental vibrational echo decay data at  $T_w = 2$  ps and (b) linear spectrum for *Ht*-M61A (dashed lines) overlaid with the best-fit vibrational echo decay and linear spectrum calculated from nonlinear response theory (solid lines) at  $1975 \text{ cm}^{-1}$ .

The linear absorption spectrum is very sensitive to the constant,  $\Delta_0$ , because its width and shape are determined by both the dynamic and the inhomogeneous contributions to the spectrum. The shapes of the vibrational echo curves and their change in shape with  $T_w$  are very sensitive to the other parameters.

The FFCF of *Ht*-M61A/Q64N is characterized by a larger  $\Delta_0$  than *Ht*-M61A, which shows that some, but not all, of the increase in linear line width shown in Figure 2 is due to an increase in inhomogeneous broadening. The first exponential terms in the FFCFs for both proteins have correlation times ( $\tau_1$ ) that are very fast (hundreds of femtoseconds). However, these components are very near the boundary of motional narrowing ( $\Delta\tau \ll 1$ ),<sup>58–61</sup> at which point  $\Delta$  and  $\tau$  cannot be determined independently.<sup>53</sup> In this nearly motionally narrowed regime, it is reasonable to believe that the relative time scale of the  $\tau_1$  values is correct, while the precise value of each  $\tau_1$  is not well-defined. To compare the relative dynamics encompassed by the first exponential terms in the FFCFs for *Ht*-M61A/Q64N and *Ht*-M61A, it is instructive to express this nearly motionally narrowed component of the FFCF as a pure dephasing time ( $T_2^*$ ), which depends on both  $\Delta_1$  and  $\tau_1$  (see section II.E). The  $T_2^*$  values for *Ht*-M61A/Q64N and *Ht*-M61A are 2.9 and 4.1 ps, respectively. That the dephasing dynamics represented by this first exponential component of  $C(t)$  are over 40% faster for *Ht*-M61A/Q64N than *Ht*-M61A is consistent with the faster vibrational echo decay shown in Figure 3. Likewise, the similar longer time scale dynamics shown for these two mutants in Figure 4 is reflected in the second exponential component of their FFCFs. Since this component of  $C(t)$  is not motionally narrowed, both the magnitude ( $\Delta_2$ ) and the correlation time ( $\tau_2$ ) are quantitatively correct and can be used to describe both the vibrational echo and the linear-IR data. The extracted FFCFs

show that the CO frequency fluctuations that occur on longer time scales (tens of picoseconds) for *Ht*-M61A/Q64N and *Ht*-M61A are characterized by very similar, but not identical,  $\Delta_2$  and  $\tau_2$  values, supporting the similar shape of the vibrational echo peak shift data in Figure 4. The fundamental dynamical difference imparted by replacing Gln64 with Asn64 is an increase in dephasing on the hundreds of femtoseconds time scale, while the CO frequency fluctuations on the tens of picoseconds time scale are virtually unchanged.

**C. Multidimensional NMR Studies: Nuclear Overhauser Effect Spectroscopy.** To determine the locations of residue 64 in *Ht*-M61A and *Ht*-M61A/Q64N relative to the heme pocket, we performed multidimensional NMR studies to complement the multidimensional IR experiments described above. Analysis of  $^1\text{H}$  NMR data focused on assigning  $^1\text{H}$  resonances for Gln64 in *Ht*-M61A and Asn64 in *Ht*-M61A/Q64N and identifying NOEs to these residues. Comparison of NMR data collected on both mutants studied above (*Ht*-M61A and *Ht*-M61A/Q64N) to published results on native *Ht*-cyt  $c_{552}$ <sup>17,50</sup> and a variant with the Q64N but not the M61A mutation (*Ht*-Q64N)<sup>22</sup> facilitated this analysis. The residue 64 conformations seen in *Ht*-cyt  $c_{552}$  and in *Ht*-Q64N are readily distinguished from each other on the basis of NOE patterns. In wild-type reduced *Ht*-cyt  $c_{552}$ , assignment of  $^1\text{H}$  NMR resonances for Gln64 was made by identification of NOEs from a Gln side chain to the Asn65 NH (which, in turn, shows NOEs to Val66). In addition, Gln64  $\epsilon$ -NH<sub>2</sub> protons have NOEs to side-chain protons of Met61 and to heme 3-CH<sub>3</sub>.<sup>22,50</sup> In *Ht*-Q64N, Asn64 is also assigned by identification of NOEs from its side chain to Asn65 NH. The Asn64  $\delta$ -NH<sub>2</sub> protons have additional NOEs to the side chains of Ile48 and Met61, heme  $\alpha$ -*meso*-H, and heme thioether-2-CH<sub>3</sub>.<sup>22</sup>

Figure 1 provides schematic representations of the heme pockets of (a) *Ht*-M61A/Q64N and (b) *Ht*-M61A. In *Ht*-M61A/Q64N, Asn64 is readily assigned by identification of NOEs from its side chain to Asn65 NH at 9.14 ppm. In addition, NOEs are observed from one or both Asn64  $\delta$ -NH<sub>2</sub> protons to Val66  $\gamma$ -CH<sub>3</sub>, Ile48  $\beta$ -H, heme  $\alpha$ -*meso*-H, and Ala61  $\beta$ -CH<sub>3</sub> but not to heme 3-CH<sub>3</sub>. The pattern of NOEs for Asn64 in *Ht*-M61A/Q64N is similar to that seen for Asn64 in *Ht*-Q64N and in *P. aeruginosa* cyt  $c_{551}$ ,<sup>22,77</sup> suggesting that Asn64 *Ht*-M61A/Q64N is oriented above the heme iron as in the typical cyt  $c_8$ 's and thus is positioned to interact with an axial CO ligand (Figure 1a). In *Ht*-M61A, Gln64 was also assigned by its characteristic NOEs to Asn65 and Val66. The Gln64  $\epsilon$ -NH<sub>2</sub> protons display NOEs to Val66  $\gamma$ -CH<sub>3</sub>'s and to heme 3-CH<sub>3</sub>, whereas NOEs to Ile48 and Ala61 are not observed. This pattern is consistent with an orientation away from the heme iron, toward the heme 3-CH<sub>3</sub> near the protein surface, as seen in the wild-type *Ht*-cyt  $c_{552}$  (Figure 1b). A summary of the multidimensional NMR results is provided in Table 2.

In addition to NOEs, chemical shifts provide an important indication of the position of an amino acid relative to the heme. Nuclei oriented above the center of the heme macrocycle, in position to interact with an axial ligand, experience a substantial upfield chemical shift via the heme ring current.<sup>78</sup> This effect has been shown to produce a readily identifiable pattern of upfield chemical shifts for the axial Met in reduced cyt  $c$ 's.<sup>79,80</sup> In both *Ht*-cyt  $c_{552}$  variants in this study, the Ala61  $\beta$ -CH<sub>3</sub> is shifted upfield as expected if its position is similar to the  $\beta$ -CH<sub>2</sub> in Met61 in the native *Ht*-cyt  $c_{552}$  ( $\delta \approx 1.5$  ppm; Table 2). In the case of an Asn64 (or Gln64) side-chain NH<sub>2</sub> proton, location above the center of the heme macrocycle (i.e., in position to interact with an axial ligand) would mean a significantly lower

**TABLE 2:**  $^1\text{H}$  NMR Data on Reduced *Ht* Cyt  $c_{552}$  Derivatives

protein	position	residue	atom(s)	$\delta$ (ppm)	relevant NOEs	refs
<i>Ht</i> cyt $c_{552}$	61	Met	$\beta$ -CH <sub>2</sub>	2.59, 0.61	heme 5-CH <sub>3</sub> , $\beta$ - <i>meso</i> -H; G52 NH	15, 50
	64	Gln	NH	7.29	M61 $\epsilon$ -CH <sub>3</sub> ; N65 NH	
	64	Gln	$\epsilon$ -NH <sub>2</sub>	6.37, 6.67	heme 3-CH <sub>3</sub> , $\alpha$ - <i>meso</i> -H; M61 $\beta$ -H, $\epsilon$ -CH <sub>3</sub>	
<i>Ht</i> -Q64N	61	Met	$\beta$ -CH <sub>2</sub>	2.78, 0.81	heme $\beta$ - <i>meso</i> -H; G52 NH	22
	64	Asn	NH	7.00	N65 NH	
	64	Asn	$\delta$ -NH <sub>2</sub>	3.18, 7.49	heme 2-CH <sub>3</sub> , $\alpha$ - <i>meso</i> -H; I48 $\gamma$ -H; M61 $\gamma$ -H, $\epsilon$ -CH <sub>3</sub> ; V66 $\gamma$ -CH <sub>3</sub>	
<i>Ht</i> -M61A	61	Ala	$\beta$ -CH <sub>3</sub>	1.48	heme 5-CH <sub>3</sub> , $\beta$ - <i>meso</i> -H	this work
	64	Gln	NH	not observed		
	64	Gln	$\epsilon$ -NH <sub>2</sub>	6.04, 6.19	heme 3-CH <sub>3</sub> ; V66 $\gamma$ -CH <sub>3</sub>	
<i>Ht</i> -M61A/Q64N	61	Ala	$\beta$ -CH <sub>3</sub>	1.53	heme $\beta$ - <i>meso</i> -H, $\gamma$ - <i>meso</i> -H; G52 NH; N64 $\delta$ -H	this work
	64	Asn	NH	7.10	N65 NH	
	64	Asn	$\delta$ -NH <sub>2</sub>	2.76, 6.85	heme $\alpha$ - <i>meso</i> -H; I48 $\beta$ -H; A61 $\beta$ -CH <sub>3</sub> ; V66 $\beta$ -CH <sub>3</sub>	

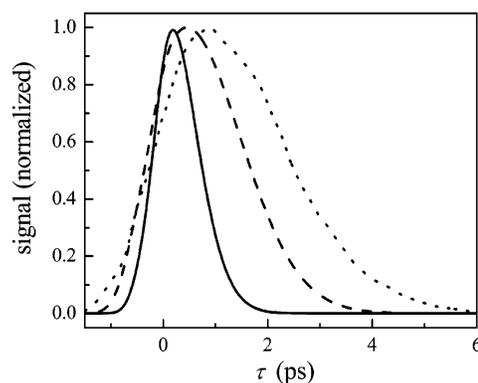
chemical shift from the expected value of  $\sim 6$ – $8$  ppm (“random-coil” chemical shifts are 6.8 and 7.6 ppm<sup>55</sup>). For example, an upfield shift of  $\sim 3$ – $5$  ppm for one Asn64  $\delta$ -NH<sub>2</sub> proton ( $\delta = 3.19$  ppm) via the heme ring current is observed in reduced *P. aeruginosa* cyt  $c_{551}$ .<sup>22,77</sup> The other Asn64  $\delta$ -NH<sub>2</sub> proton, which is oriented away from the axial ligand, has a more typical shift of 7.49 ppm.<sup>81</sup> This pattern is seen for the Asn64  $\delta$ -NH<sub>2</sub> protons in other proteins in the cyt  $c_8$  family: *Nitrosomonas europaea* cyt  $c_{552}$  (3.35, 7.11 ppm),<sup>82</sup> *Pseudomonas stutzeri* cyt  $c_{551}$  (3.20, 7.02 ppm),<sup>20</sup> and *Pseudomonas stutzeri* Zobell cyt  $c_{551}$  (3.11, 6.95 ppm).<sup>21</sup> In *Ht*-M61A/Q64N, the Asn64  $\delta$ -NH<sub>2</sub> protons have chemical shifts (3.18, 7.49 ppm) that are consistent with an orientation similar to that seen in the typical cyt  $c_8$ 's.<sup>22</sup> In reduced *Ht*-cyt  $c_{552}$ , in contrast, the chemical shifts for the Gln64  $\epsilon$ -NH<sub>2</sub> protons (6.37, 6.67 ppm) are *not* consistent with a side-chain position near the heme iron or its axial ligand.<sup>17,22,23,50</sup> In the current study, analogous chemical shift patterns are seen for Asn64  $\delta$ -NH<sub>2</sub> in *Ht*-M61A/Q64N ( $\delta = 2.76$ , 6.85 ppm) and Gln64  $\epsilon$ -NH<sub>2</sub> in *Ht*-M64A ( $\delta = 6.04$ , 6.19 ppm). These NMR data provide further support for the hypothesis that Asn64 in *Ht*-M61A/Q64N is positioned to interact with the axial CO, whereas Gln64 in *Ht*-M61A is not.

In light of the structural characterization of *Ht*-M61A and *Ht*-M61A/Q64N by NMR, the comparison of these *Ht*-cyt  $c_{552}$  mutants to MbCO and its variants can now be elaborated. As described above, the linear-IR spectra of *Ht*-M61A and *Ht*-M61A/Q64N (Figure 2) appeared to correspond structurally to the MbCO A<sub>0</sub> and A<sub>1</sub> or A<sub>3</sub> conformational substates, respectively. The NMR results clearly indicate that the structural analogy is valid: *Ht*-M61A and the A<sub>0</sub> substate are characterized by residue 64 directed out of the heme pocket, while this residue is positioned within the pocket above the heme iron in *Ht*-M61A/Q64N and the A<sub>1</sub> and A<sub>3</sub> substates. In general, these data indicate that the inclusion of a polar hydrogen-bond-donating residue above the heme ring has a noticeable effect on the active site dynamics. For comparison to the *Ht*-cyt  $c_{552}$  mutant dephasing dynamics shown in Figure 3, the spectrally resolved vibrational echo decays for MbCO at the A<sub>1</sub> (dashed curve) and A<sub>3</sub> (solid curve) substates are presented in Figure 5. Due to spectral overlap of all three spectroscopic lines in MbCO, the echo data for these substates are complicated by accidental degeneracy beats (ADB).<sup>83</sup> Fortunately, the FFCF extraction procedure described in section II.E allows the vibrational echo decays for each substate to be recalculated, as shown by the A<sub>1</sub> and A<sub>3</sub> echo decays in Figure 6, without the influence of the other states. The low intensity of the A<sub>0</sub> substate precludes the acquisition of data at this frequency; however, a MbCO mutant in which H64 has been replaced by a valine (H64V) has been shown to represent the CO vibrational dynamics that correspond to this substate.<sup>51</sup> The vibrational echo decay for H64V (dotted curve), representing the A<sub>0</sub> substate decay, is overlaid in Figure

6 with the decays for the A<sub>1</sub> and A<sub>3</sub> substates (all decays shown at  $T_w = 0.5$  ps).

It is apparent in Figure 6 that the dephasing of the A<sub>3</sub> substate is faster than that of the A<sub>1</sub> substate,<sup>39,56,84</sup> which in turn dephases faster than the A<sub>0</sub> substate. The fundamental difference between the A<sub>1</sub> and the A<sub>3</sub> structures lies in the rotation of the singly protonated imidazole ring on His64. The N <sub>$\epsilon$</sub> -H proton and N <sub>$\delta$</sub>  of this imidazole ring are equidistant from the CO ligand in the A<sub>1</sub> substate, whereas the N <sub>$\epsilon$</sub> -H proton is directed toward the CO ligand in the A<sub>3</sub> substate.<sup>39,56</sup> In the A<sub>3</sub> geometry, the direction of the hydrogen-bond-donating group (N <sub>$\epsilon$</sub> -H proton) toward the heme-bound CO could provide an additional source of dephasing and generate the faster vibrational echo decay in Figure 6. It is important to recognize that *Ht*-cyt  $c_{552}$  and MbCO are different proteins within the general category of heme proteins, and their overall dynamic ranges are quite different (Figures 3 and 5). Nonetheless, the relative magnitude of the change of the dephasing rate from *Ht*-M61A to *Ht*-M61A/Q64N is more similar to that of the MbCO A<sub>0</sub> to the A<sub>1</sub> substate than that of A<sub>0</sub> to A<sub>3</sub>. While this is admittedly pushing the limitations of the analogy, the similarity of *Ht*-M61A/Q64N dynamics (relative to *Ht*-M61A) to the A<sub>1</sub> substate could suggest that the hydrogen-bonding amine group on Asn64 is rotated away from the CO ligand to some degree. To unambiguously identify the atomic displacements responsible for the vibrational dynamics measured in these experiments necessitates calculation of the vibrational echo data from molecular dynamics simulations.

That the dynamical trends of *Ht*-M61A and *Ht*-M61A/Q64N are similar to those of the conformational substates of MbCO is an intriguing result. In the case of MbCO, His64 has been implicated in the physiologically crucial differentiation between CO and O<sub>2</sub> binding to the heme group of myoglobin and



**Figure 6.** Spectrally resolved vibrational echo decays at  $T_w = 0.5$  ps for CO bound to MbCO at the A<sub>3</sub> (1938  $\text{cm}^{-1}$ , solid line) and A<sub>1</sub> (1944  $\text{cm}^{-1}$ , dashed line) conformational substates. The vibrational echo decay at  $T_w = 0.5$  ps for the H64V mutant is overlaid (1968  $\text{cm}^{-1}$ , dotted line) and represents the echo decay from the A<sub>0</sub> substate in MbCO.

hemoglobin.<sup>67,85–87</sup> The sensitivity of the active site dynamics for MbCO to the geometry and dynamics of residue 64 suggests that an amino acid residue positioned to interact with heme axial ligands also has a profound effect on the functionality of the active site of cyt c's. This hypothesis has been recently supported by dynamical studies on *Ht*-cyt c<sub>552</sub> and *Ht*-M61A/Q64N.<sup>17,22</sup> It is plausible that the presence of Gln64 in the active site of *Ht*-cyt c<sub>552</sub> instead of the Asn64 found in most other cyt c<sub>8</sub>'s affords a unique physiological function that is beneficial to that species. The current study reveals a dynamical difference imparted by Gln64 on the picosecond time scale and shows that, in addition to affecting slower structural motions, single site mutations in naturally occurring proteins can also influence dynamical processes. Fast structural fluctuations can be the precursors to slower time scale structural changes.

#### IV. Concluding Remarks

A protein's function is defined by its structural architecture and the evolution of that architecture with time. To understand, alter, or mimic the function of a protein or enzyme, it is necessary to understand the roles that specific residue motions play in the determining the physiological reactivity of the active site. The multidimensional IR experiments presented here revealed that the rate of dephasing (picosecond time scale) of the heme-bound CO in *Ht*-M61A/Q64N is significantly faster than that for *Ht*-M61A, while the rate of spectral diffusion (tens of picoseconds time scale) is nearly identical for the two mutants. This implies that the crucial residues surrounding the active site of a protein or enzyme could be optimized to satisfy the specific needs of an organism without significantly altering the longer time scale structural dynamics that typically involve movements of larger domains. Multidimensional NMR experiments provided data that were complementary to the multidimensional IR studies and determined that the geometry of residue 64 in the two mutants corresponded to a non-hydrogen-bonding and hydrogen-bonding interaction for *Ht*-M61A and *Ht*-M61A/Q64N, respectively. We conclude that the faster dynamics on the picosecond time scale measured in *Ht*-M61A/Q64N are due to the geometry of Asn64, which is a hydrogen bond donor that localizes above the heme-bound CO. A similar interaction between residue 64 and the CO ligand has been observed for MbCO and its H64V variant. These examples suggest a general trend toward rapid active site dynamics in the presence of a hydrogen bond donor and represent a mechanism by which an organism might impart a unique selectivity to substrate binding or reactivity at the active site of a protein or enzyme.

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**Supporting Information Available:** Experimental vibrational echo decay data and the linear spectrum for *Ht*-M61A and *Ht*-M61A/Q64N overlaid with the best-fit vibrational echo decay and linear spectrum calculated from nonlinear response theory. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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